Quantum Confinement Controls Photocatalysis: A Free Energy Analysis for Photocatalytic Proton Reduction at CdSe Nanocrystals

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ABSTRACT The ability to adjust the mechanical, optical, magnetic, electric, and chemical properties of materials via the quantum confinement effect is well-understood. Here, we provide the first quantitative analysis of quantum-size-controlled photocatalytic H2 evolution at the semiconductor—solution interface. Specifically, it is found that the hydrogen evolution rate from illuminated suspended CdSe quantum dots in aqueous sodium sulfite solution depends on nanocrystal size. Photoelectrochemical measurements on CdSe nanocrystal films reveal that the observed reactivity is controlled by the free energy change of the system, as determined by the proton reduction potential and the quasi-Fermi energy of the dots. The corresponding free energy change can be fitted to the photocatalytic activity using a modified Butler—Volmer equation for reaction kinetics. These findings establish a quantitative experimental basis for quantum-confinement-controlled proton reduction with semiconductor nanocrystals. Electrochemical data further indicate that proton reduction occurs at cadmium sites on the dots, and that charge separation in these nanocrystals is controlled by surface effects, not by space charge layers.

KEYWORDS: quantum confinement · quantum dot · photolysis · water splitting · solar fuel · Butler—Volmer · space charge layer

Semiconductor and metal nanocrystals with sizes below the Bohr exciton radius exhibit quantum confinement effects, which modify the optical,1–4 electrical,5 magnetic,6,7 and chemical8,9 properties of materials. These tunable properties make semiconductor nanocrystals especially interesting for solar energy conversion applications and as photocatalysts.10–22 According to Marcus theory, quantum-confined systems support higher rates of interfacial charge transfer, due to the increased thermodynamic energy of the confined electron–hole pairs.23,24 For solid—solid interfaces, quantum-controlled electron injection has been experimentally confirmed for CdTe/CdSe core–shell and nanorod structures25,26 and for CdSe and PbS dots immobilized on TiO2.27–30 As Kamat’s group showed recently,27–32 the dependence of the interfacial charge transfer rates on the energetics of these systems can be quantitatively understood using Marcus theory. For solid—liquid interfaces, a comparable theoretical analysis is still lacking, even though quantum-confinement-controlled redox reactions have been observed for a number of systems. For electron transfer between CdSe dots (2.3–6.3 nm in diameter) and methylviologen cations, Wachtveitl and co-workers noted an exponential dependence of the electron transfer rate on the free energy change.33 For CdSe dots with adsorbed Re-bipyridyl complexes,34 or covalently linked fullerene derivatives,31 photochemical electron transfer rates were found to increase with decreasing size of the dots. Several instances of a quantum effect in photocatalytic reactions were also reported but only qualitatively. For example, WO3 nanocrystal suspensions showed a size dependence for the decomposition of benzene35 and for photocatalytic water oxidation.36 Cadmium sulfide nanoparticles were reported to exhibit a size-dependent activity for...
photoreduction of aromatic ketones and for methanol dehydrogenation. Cadmium selenide nanoribbons with a 2.7 eV expanded band gap were found to photocatalytically produce hydrogen from irradiated Na$_2$SO$_3$/NaSH solution, whereas bulk CdSe ($E_g = 1.7$ eV) was not active. We recently reported that the hydrogen evolution rates from illuminated CdSe quantum dot (QD) suspensions in Na$_2$SO$_3$ solution showed a logarithmic dependence on the semiconductor band gap. We have now been able to measure the precise energetics of the dots with electrochemical techniques and perform a thorough analysis of the observed reactivity trend. The results show that the kinetics of hydrogen evolution in this system can be understood quantitatively using a modified Butler–Volmer model, which relates the kinetic activation energy with the thermodynamic free energy change for proton reduction. These findings establish an experimental basis for quantum-size-controlled photocatalytic hydrogen evolution with nanocrystals. The data further show that proton reduction proceeds on the cadmium sites on the CdSe QD surface, and that charge transport in CdSe QDs occurs by diffusion, not by drift, as postulated by Kronik, Gratzel, Hagfeld, and Memming. These results have an impact on the use of quantum-size-confined semiconductors for solar energy conversion and on the understanding of reaction kinetics in nanostructured photocatalysts and photoelectrodes.

RESULTS AND DISCUSSION

Monodisperse CdSe quantum dots with a diameter ranging from 1.8 to 6.0 nm were synthesized using 2-mercaptoethanol as a short-chain surface-passivating agent (Figure 1A,B) based on a known procedure by Rogach et al., followed by size-selective precipitation with 2-propanol. According to HR-TEM (Figure 1C) and powder XRD (Supporting Information Figure S1) data, these QDs are single crystals of the cubic zinc blende structure type. The preference of the zinc blende over the wurtzite structure type is well-known for CdSe quantum dots. Optical properties of the QD powders (Figure 1A) are size-dependent, which is a direct outcome of the increase in their band gaps. Optical band gaps for these QDs range from 2.04 to 2.90 eV (absorption onsets between 607 and 427 nm) for sizes between 4.1 and 2.0 nm according to their UV/vis absorption spectra (Figure 1D), while bulk CdSe has a band gap of 1.76 eV (704 nm). This widening of the optical band gap as size decreases into the quantum confinement regime has been modeled and observed for many materials and is consistent with many previous studies on CdSe QDs.

To determine the photocatalytic activity of these dots, particle suspensions were irradiated with light (>330 nm, 400 mW/cm$^2$) in the presence of 0.10 M aqueous Na$_2$SO$_3$ as a sacrificial agent. Sodium sulfite was selected instead of sodium hydrosulfide because the latter was found to induce precipitation of the CdSe QDs via cross-linking. Sulfite is less active than sulfide but does protect CdSe against photocorrosion by reacting with the photogenerated holes, as shown in the inset of Figure 2B. We have previously reported that, under these conditions, all QD samples, except for the largest particles and bulk CdSe, evolved hydrogen linearly over time (Figure 2A). Interestingly, a lag occurs between the beginning of illumination and the beginning of H$_2$ evolution. This lag ranges between...
Figure 2. (A) Experimental H₂ evolution from CdSe QDs in 0.10 M Na₂SO₃ aqueous solution at pH 9.2. Inset: Photos showing visible color change of CdSe QD solution before (left) and during (right) irradiation. (B) Relative H₂ evolution rates as normalized with regard to catalyst amount and absorbed photons plotted vs QD size. Inset: Schematic diagram of hydrogen evolution at CdSe QDs in the presence of sodium sulfite.

0.5 and 3 h and roughly increases with inverse particle size. It is accompanied by the formation of a brown color (inset of Figure 2A), which disappears upon exposing the dots to air. This suggests that the color change corresponds to the in situ formation of Cd⁰ on the dot surface, which might be necessary for the catalytic hydrogen evolution. The longer lag times for the smaller CdSe dots could be interpreted as a result of the greater cadmium nucleation barrier. Attempts to directly observe the formed cadmium particles with electron microscopy failed due to the small size and the air sensitivity of the formed cadmium. We note, however, that the formation of metal particles during photocatalysis has previously been reported for ZnS photocatalysts. Correcting the H₂ evolution rates in Figure 2A for catalyst mass and absorbed photons (for details see Supporting Information Figure S2) and plotting them versus size gives the trend in Figure 2B. A monotonic increase in hydrogen evolution rate with decreasing particle size is evident. This suggests that hydrogen evolution activity of these QDs is regulated by the energetics of the charge carriers in the particles.

In order to quantitatively understand the relationship between QD size, QD energetics, and their photocatalytic activity, the energetics of the QDs were studied in detail with electrochemistry and photoelectrochemistry. Electrodes of quantum dot films (1.0 × 1.0 cm²) were prepared by drop-coating QD solutions on F/SnO₂ substrates and drying in the dark overnight. The films were immersed in an aqueous electrolyte containing 1.0 M KCl and 0.10 M Na₂SO₃ at a pH of 9.2 in all of the measurements. Sulfite was used as the sacrificial donor over sulfide because, as a potential determining ion, sulfide can fix the Fermi level of CdSe. The use of sulfite also ensured compatibility with the photocatalytic hydrogen evolution experiments. Photocurrent scans were performed cathodically from 0.4 to −1.7 V vs NHE under illumination with chopped light from a xenon arc lamp. The results are shown in Figure 3 for various sizes.

In general, photocurrents were small for bulk CdSe (60 μA/cm²) and for small QDs below 2.6 nm (10 μA/cm²) but larger for large QDs above 2.6 nm (300 μA/cm²). This trend can be explained by several competing factors, which includes an increase in surface area in going from bulk CdSe to large QDs and a decrease of the photon absorption in going from large to smaller QDs. Besides, the non-annealed QD films also have a large electrical resistance since electron transport has to occur by thermally activated hopping and/or electron tunneling between dots. Smaller QDs lead to more junctions between the dots, thus a higher resistance is expected.

The most interesting aspect of the scans is that QDs with diameters above 2.6 nm produced large anodic photocurrents, whereas QDs with diameters below 2.6 nm produced small cathodic photocurrents, as shown in Figure 3A,B. For large dots (>2.6 nm), the size of the anodic current was a strong function of the applied potential and occurred only positive of an onset potential E₀, that was characteristic for each CdSe film (Figure 3A). For instance, E₀ varies from −0.51 V for bulk CdSe to −0.75 V for 4.0 nm dots and −1.13 V for 2.8 nm dots. On the contrary, small QDs (<2.6 nm) produced only a small (10 μA/cm²) cathodic photocurrent, whose magnitude was nearly independent of the applied potential or QD diameter (Figure 3B). Interestingly, all of the cathodic photocurrents showed a similar photo-onset E₀, near −0.3 V. In addition, a small anodic photocurrent at potentials positive of −0.1 V was observed, but this signal originated from the F/SnO₂ electrode background, not from the QD film (see bottom trace in Figure 3B). The observed photocurrent inversion in going from large to small CdSe QDs is not unprecedented in the literature. It has been previously observed by Kronik and co-workers for CdSe dot films from before to after hydrochloric acid etching and been assigned to the charge trapping at surface sites. Our results and the following discussion support this interpretation.
As an n-type material, bulk CdSe normally forms a depletion layer when in contact with a suitable redox couple in the liquid phase (e.g., H⁺/H₂). However, the QDs investigated here are too small to allow a full space charge layer to develop. Assuming a donor concentration of \( n_D = 1.5 \times 10^{17} \) cm\(^{-3}\), a dielectric constant of \( \varepsilon_r = 9 \) for CdSe, and a full depletion case, the space charge layer is 80 nm wide; 20 times larger than the diameter of the dots. Under these conditions, the potential barrier within the dots is limited by the nanocrystal diameter and does not exceed 3.9 meV for a 2.0 nm dot or 17 meV for a 4.2 nm dot, as calculated using the method by Albery. Such a barrier is not effective for electrons at room temperature (thermal energy \( kT = 26 \) meV), which means that, instead of surface band bending, other factors determine the charge separation at the QD–liquid interface. We attribute the observed inversion of the photocurrent to the preferential trapping of electrons at Cd\(^{2+}\) surface states, as shown in Figure 3C,D. For the large QDs (>2.6 nm), these Cd\(^{2+}\) surface states are not accessible because their energy is above the quasi-Fermi energy of the photogenerated electrons \( E_{F,n} \). As a result, an anodic photocurrent is observed due to the sulfite oxidation by photogenerated holes and the extraction of photoelectrons at the back electrode when the applied potential is positive of \( E_{F,n} \). Thus, the photo-onset \( E_{F,n} \) is a measure of the Fermi energy of the electrons \( E_{F,n} \) in the dots. In the small QDs (<2.6 nm), \( E_{F,n} \) has been raised due to quantum confinement effect. Consequently, these photogenerated electrons become preferentially trapped at surface Cd\(^{2+}\) ions to produce Cd\(^0\). A cathodic photocurrent is observed, due to the extraction of photogenerated holes by the back electrode when the applied potential at the electrode is negative of the quasi-Fermi level for holes \( E_{F,h} \). Due to the low mobility of the holes, the cathodic photocurrents in the small dots are much lower than the anodic photocurrents in the large dots. Importantly, the onset potential for the cathodic current in the small dots is no longer a good measure of the quasi-Fermi level for electrons \( E_{F,n} \). As it turns out, \( E_{F,n} \) can be estimated from electrochemical scans in the dark instead. Three reduction features can be discerned in the dark scans, which are labeled as I, II, and III (Figure 4).
Features I and II are associated with each other and shift cathodically from −0.73 and −1.05 V for dots of 4.2 nm, to −1.06 and −1.35 V for dots of 3.0 nm, and −1.14 and −1.37 V for dots of 2.8 nm (Figure 4, Table 1, and Figure S3). These reduction peaks do not lead to any visible gas evolution but induce a visible brown coloration of the films. The same color was also observed during irradiation of the dots in the presence of Na2SO3 (Figure 2A inset). Thus, reduction peaks I and II are attributed to the reduction of surface and lattice Cd2⁺ ions, respectively. The observed reduction peak potentials compare well to literature values.

For example, 6.5 nm TOPO-coated CdSe dots undergo Cd2⁺ reduction at −1.06 V vs NHE in acetonitrile,\(^7\) and for 2.8 nm thioglycolic acid-capped CdTe dots, Cd2⁺ reduction is observed at −1.40 V vs NHE.\(^8\) Adding electrons to lattice Cd2⁺ sites is equivalent to adding electrons to the CdSe conduction band; therefore, the reduction peak potential is a good indication of \(E_{\text{CB}}\) in the CdSe dots and is listed in Table 1, together with other parameters. Indeed, for QDs with diameters larger than 2.6 nm, Cd2⁺ reduction features in Figures 4 and S3 occurred slightly negative of the photo-onset potentials \(E_{\text{ph}} = E_{\text{fb}}\) in Figure 3A. For example, in 4.2 nm dots, \(E_{\text{ph}} = −0.72\) V and \(E(\text{Cd}^{2+}/\text{Cd}) = −1.05\) V, and for 2.9 nm dots, \(E_{\text{ph}} = −1.07\) V and \(E(\text{Cd}^{2+}/\text{Cd}) = −1.33\) V. The observed \(E_{\text{CB}}−E_{\text{fb}}\) separation of 0.25–0.30 V indicates that the QDs are weakly n-doped. For QDs with diameters below 2.6 nm, the Cd2⁺ reduction peak shifts further cathodically and merges with feature III near −1.6 V, only observable as a shoulder (Figures 4 and S3). Therefore, for the smallest dots (<2.6 nm), \(E_{\text{CB}}\) is determined from the shoulder potential. As \(E_{\text{fb}}\) in these dots cannot be obtained from the cathodic photocurrent onsets, \(E_{\text{fb}}\) values are estimated using the measured \(E_{\text{CB}}\) to be negative of −1.13 V (\(E_{\text{fb}}\) of the closest 2.8 nm QD) and at least 0.059 V positive of the corresponding \(E_{\text{CB}}\) (Table 1).

Cathodic feature III in the dark scan (Figure 4) occurs near −1.6 V. This feature is associated with the reduction of protons, as evident from the observed \(\text{H}_2\) gas evolution at the electrode. A Tafel plot in Figure S4 presents the linear regime of the logarithmic current density \(\log(i)\) versus the applied potential curves for proton reduction over various QD films.

\[
\log(i) = -\frac{\alpha F}{2.303RT} E + \log(i_0) \tag{1}
\]

As indicated by eq 1 for the Tafel plot, extrapolating the linear fit to zero applied potential yields the exchange current density \(i_0\) from the intercept. Values for the different QD films fall into a narrow interval between

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**TABLE 1. Summary of Energetic Properties for CdSe QDs**

<table>
<thead>
<tr>
<th>diameter/nm</th>
<th>(E_f/\text{eV})</th>
<th>(E_f^*/\text{V})</th>
<th>(E_{\text{fb}}/\text{V})</th>
<th>(E_{\text{CB}}/\text{V})</th>
<th>(E(\text{Cd}^{2+}/\text{Cd})_\text{surf}/\text{eV})</th>
<th>normalized (\text{H}_2) evolution rate (mole H2/mole CdSe)</th>
</tr>
</thead>
<tbody>
<tr>
<td>bulk</td>
<td>1.74</td>
<td>−0.51</td>
<td>NA</td>
<td>−0.85</td>
<td>0.89</td>
<td>−0.67</td>
</tr>
<tr>
<td>4.2</td>
<td>1.97</td>
<td>−0.72</td>
<td>−0.72</td>
<td>NA</td>
<td>−1.05</td>
<td>0.92</td>
</tr>
<tr>
<td>4.0</td>
<td>2.05</td>
<td>−0.75</td>
<td>−0.75</td>
<td>−1.05</td>
<td>1.00</td>
<td>0.75</td>
</tr>
<tr>
<td>3.0</td>
<td>2.39</td>
<td>−1.03</td>
<td>−1.03</td>
<td>−1.35</td>
<td>1.04</td>
<td>−1.06</td>
</tr>
<tr>
<td>2.9</td>
<td>2.43</td>
<td>−1.07</td>
<td>−1.07</td>
<td>1.33</td>
<td>1.10</td>
<td>−1.09</td>
</tr>
<tr>
<td>2.8</td>
<td>2.47</td>
<td>−1.13</td>
<td>−1.13</td>
<td>1.37</td>
<td>1.10</td>
<td>−1.14</td>
</tr>
<tr>
<td>2.5</td>
<td>2.58</td>
<td>−0.24</td>
<td>−1.13</td>
<td>−1.13</td>
<td>−0.24</td>
<td>−1.37</td>
</tr>
<tr>
<td>2.4</td>
<td>2.68</td>
<td>−0.27</td>
<td>−1.13</td>
<td>−1.38</td>
<td>−0.27</td>
<td>−1.44</td>
</tr>
<tr>
<td>2.3</td>
<td>2.73</td>
<td>−0.28</td>
<td>−1.13</td>
<td>−1.48</td>
<td>−0.28</td>
<td>−1.54</td>
</tr>
<tr>
<td>2.0</td>
<td>2.88</td>
<td>−0.33</td>
<td>−1.13</td>
<td>−1.59</td>
<td>−0.33</td>
<td>−1.65</td>
</tr>
</tbody>
</table>

\(^a\) All potentials were measured at pH 9.2 with 0.10 M Na2SO3 as the electron donor and are reported vs NHE. \(^b\) \(E_{\text{fb}}\) is designated as the onset of the photocurrent in the photoelectrochemical scans in Figure 3A. \(^c\) \(E_{\text{ph}}\) is determined from the photocurrent of the anodic photocurrent for dots above 2.6 nm in Figure 3A. \(^d\) \(E_{\text{fb}}\) is determined from the photo-onset of the cathodic photocurrent for dots below 2.6 nm in Figure 3B. \(^e\) \(E_{\text{CB}}\) is assigned from the peak potential of the lattice Cd2⁺ reduction feature II in Figure 4 and Figure S3. \(^f\) \(E_{\text{CB}}\) is calculated using \(E_f = E_{\text{ph}} + E(\text{Cd}^{2+}/\text{Cd})_\text{surf}\) measured from the surface Cd2⁺ reduction potential (feature I in Figure 4 and Figure S3). \(^g\) \(E_{\text{ph}}\) range estimated from \(E_{\text{ph}}\) of 2.8 nm QD to 0.059 V below measured \(E_{\text{ph}}\) values. \(^h\) Unit: moles H2/(h×mole CdSe×overlap integral area).
The Tafel plots in Figure S4 establish cadmium metal overpotentials, 10^{-10} to 10^{-12} A/cm² for Cd, 10^{-12} A/cm² for Zn, and 10^{-12} A/cm² for Hg. This supports the previous hypothesis that electrochemical proton reduction occurs on the Cd²⁺ sites present in CdSe QDs.

The energetics data for all CdSe QDs are summarized in Table 1 and in Figure 5A. Several features are notable. As predicted by quantum size effects, the band energy more reducing and the VB more oxidizing with diminishing QD size. However, the observed shift in $E_{CB}$ from 0.85 to -1.65 V is much more prominent than the shift in $E_{VB}$ from 0.89 to 1.23 V. This difference is due to the effect of different effective masses of electrons $m_e^*$ and holes $m_h^*$. As shown in eq 2, the quantum-confined energy $E_{qh}$ for electrons and holes is a direct function of their effective mass.

$$E_{qh} = \frac{\hbar^2 \pi^2}{2m_{qh}^* a^2}$$

Here, $a$ is the particle diameter, and all other variables have their common meanings. Given that $m_e^*(CdSe) = 0.13 m_e$ and $m_h^*(CdSe) = 0.45 m_h$, a stronger dependence of $E_{CB}$ on the size of the QDs than that of $E_{VB}$ is expected. This can be verified by more rigid computations.

For example, directly plotting the conduction band edge as a function of size gives a trendline of $E_{CB} \propto a^{-0.61}$ with a fitted exponent of $-0.61 \pm 0.04$ (see Figure S5). This agrees well with the conduction band shift measured by X-ray absorption spectroscopy (yields an exponent of $-0.6 \pm 0.1$) and computed by tight-binding approximation or charge-patching method (yields exponents of -1.0 or -0.8, respectively). Moreover, as can be seen in Figure 5A, the quasi-Fermi level of electrons $E_{F,n}$ in general 250–300 mV below the conduction band edge $E_{CB}$ and shifts along with $E_{CB}$ as size decreases.

In Figure 5B, photocatalytic hydrogen evolution rates are plotted in the logarithmic form together with quasi-Fermi energies versus QD diameters. The correlation between the logarithmic rates and the $E_{F,n}$ values is evident. Both increase as the size of the dots decreases.

In the following, we show that this correlation can be understood in terms of a free-energy-controlled interfacial charge transfer. We assume that, under the photocatalytic conditions (small reaction rate and rapid stirring), mass transport of the reactants to the surface of the catalyst is not rate-limiting. Instead, the reaction rate is controlled by the solid–liquid interfacial kinetics. As sketched in Figure 6, there are two half-reactions at the interface: the oxidation of sulfite and the reduction of protons. The driving force for the oxidation reaction $\Delta G_{ox}$ is given by the difference between $E_{F,n}$ and $E[\text{SO}_3^-/\text{SO}_4^{2-}]$ and the driving force for the reduction reaction $\Delta G_{red}$ by the difference between $E_{F,n}$ and $E[H^+/H_2]$. Compared to $\Delta G_{red}$, $\Delta G_{ox}$ for sulfite oxidation is large and relatively constant across the series of QDs. On the other hand, $\Delta G_{red}$ is small and strongly affected by the QD size variation as seen in Figure 5A. Thus, the proton reduction kinetics are considered to be the rate-limiting factor for the photocatalytic hydrogen evolution over CdSe QDs at pH 9.2.

The Tafel plots in Figure S4 establish cadmium metal as the site for proton reduction. Therefore, the cathodic electrochemical proton reduction overpotential over CdSe QDs.

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Figure 7. (A) Butler–Volmer diagram. As the QD size diminishes, the Fermi energy of electrons shifts up from $E^*$ to $E$, causing an increase of $(1-\alpha)F(E-E^*)$ in the free energy change $\Delta G$ as well as a decrease of $\alpha F(E-E^*)$ in the kinetic activation energy $\Delta G_k$. The value of the electron transfer coefficient $\alpha$ depends on the shape of the energy surface of the reactants and the products. (B) Weighted least square linear fit of the experimental data (from Table 1 and Figure 2B). Error bars for the four smallest dots are estimated to reach from $E_{F,n}$ for the 2.8 nm QD to 0.059 V below the measured $E_{C,B}$ values and are taken into the weighted least square linear fit. The lowest point is excluded from the fit (dashed line) because the rate is close to the experimental error ($\pm 10^{-4}$) of the $H_2$ rate measurement.

current density can be described by the Butler–Volmer equation for electrode kinetics as shown below:

$$j = Fk_0[H^+]e^{-\alpha F(E-E^*)/RT} - Fk_0[H_2]e^{\alpha F(E-E^*)/RT}$$

(3)

Here, $[H^+]$ is the concentration of the electron acceptor in the liquid (same for all QDs solutions as determined by pH) and $[H_2]$ the concentration of the electron donor; $k_0$ is the standard rate constant, and $F(E-E^*)$ is the driving force for the reaction, with $F$ as the Faraday constant, $E$ as the Fermi energy of electrons in the QDs $E_{F,n}$ and $E^*$ equivalent to the Nernst potential for the redox couple $E[H^+/H_2]$ in this case; $\alpha$ is the charge transfer coefficient as explained in Figure 7A, which determines the fraction of the increase in thermodynamic free energy change $\Delta G_{\text{red}} - \Delta G_{\text{red}^*} = F(E-E^*)$ that gets carried onto the reduction in kinetic activation energy for the forward reaction, as given by $\Delta G_k = \Delta G_{\text{red}} - \alpha F(E-E^*)$. Importantly, the second term for the back reaction (hydrogen oxidation) in eq 3 can be ignored because the hydrogen concentration $[H_2]$ is small compared to the proton concentration $[H^+]$ even at a mild basic pH of 9.2. Converting current density to hydrogen evolution rate by taking into consideration the stoichiometry (two electrons are required to make one $H_2$ molecule), one obtains eq 4 and its logarithmic form, eq 5, to relate the hydrogen evolution rate to the free energy change $\Delta G_{\text{red}} = F(E_{F,n} - E[H^+/H_2])$.

$$V_{H_2} = V_{\text{net}} \approx V_{\text{forward}} = \frac{j}{2eN_A} = \frac{1}{2} k_0[H^+]e^{-\alpha F(E-E^*)/RT}$$

(4)

$$\log(V_{H_2}) = -\frac{\alpha F(E_{F,n} - E[H^+/H_2])}{RT} + \log\left(\frac{1}{2} k_0[H^+]\right) = -\frac{\alpha F\Delta E}{RT} + \log\left(\frac{1}{2} k_0[H^+]\right)$$

(5)

A plot of eq 5 with experimental $H_2$ evolution rate $\log(V_{H_2})$ versus $\Delta E = E_{F,n} - E[H^+/H_2]$ is shown in Figure 7B. The model fits the experimental $H_2$ rate data well, except for the data point with the lowest rate. Here, the experimental error is large and the back reaction (hydrogen oxidation) can no longer be ignored. The fitted slope $-\alpha F/2.303RT$ is $-2.8$, giving an electron transfer coefficient $\alpha = 0.17 - 0.19$, depending on whether one uses the temperature from the electrochemical measurements (25 °C) or from the homogeneous irradiation experiments (65 °C). Literature values for $\alpha$ range between 0.15 and 0.5, depending on the electrode material.84 - 86 Note that values for $\alpha$ from Figure 7B are similar to $\alpha = 0.21 - 0.25$ from the Tafel plots in Figure S4. The match is remarkable considering the different processes involved in the photocatalytic experiments (photogenerated electrons diffuse to the QD surface to reduce protons, whereas the photogenerated holes get captured by $SO_3^{2-}$) and in the dark electrochemical experiments (electrode injects electrons through QD films and reduces protons at the film–liquid interface). This further confirms that the interfacial proton reduction kinetics are the limiting factor in the photocatalytic hydrogen evolution reaction at CdSe QDs.

CONCLUSIONS

In summary, we provide the first quantitative analysis on the size dependence of photocatalytic hydrogen evolution at the solid–liquid interface. The observed behavior is analogous to size-controlled charge transfer across solid–solid interfaces, as reported by Kamat’s group.27 - 29 The effect can be understood by relating the electrochemical electron potential in the photoexcited dots to the activation energy for proton reduction using a modified Butler–Volmer equation. This confirms that the reactivity of the dots is controlled by their free energy. Furthermore, this work shows that the observed photocurrent inversion from anodic to cathodic in the small dots can
be understood in terms of electron trapping at the surface sites, and that proton reduction occurs on Cd(0) surface sites, as confirmed by the Tafel behavior of the proton reduction current. These results are relevant to the understanding of charge transfer at nanostructured photoelectrodes and to the engineering of advanced devices for solar energy conversion.

METHODS

Chemicals. Cadmium perchlorate (95%, Alfa Aesar), selenium powder (99.5%, Acros Organics), sodium borohydride (98%, Strem Chemicals), 2-mercaptoethanol (98%, Aldrich), sodium sulfite (98%, Merck), and potassium chloride (99.6%, Fisher Scientific) were used as received. Water was purified to a resistivity of >18 MΩ by a Nanopure II system.

Aqueous Synthesis of CdSe QDs. An aqueous synthesis developed by Rogach et al.48 was adopted in this study for the preparation of CdSe QDs with 2-mercaptoethanol as the surface capping agent. After obtaining the crude product, a size refinement was performed using size-selective precipitation with 2-propanol as the nonsolvent. The size and the monodispersity of the QDs were confirmed by transmission electron microscopy (TEM), X-ray diffraction (XRD), as well as the optical band gap of the QDs using a correlation published by Schooss et al.43,44

Physical Property Characterizations. Powder XRD spectra of CdSe QDs were obtained on a Scintag XDS-2000 diffractometer with a wavelength of 0.154 nm (Cu Kα line), a tube slit divergence of 2.0 mm, a column scatter of 0.5 mm, and a receiving width of 0.2 mm. TEM samples were prepared by dipping carbon-coated copper grids into aqueous dispersions of CdSe QDs, followed by rinsing with water and air-drying. Bright-field high-resolution transmission electron microscopy (HR-TEM) images were captured on a JEOL JEM-2500 SE microscope with an accelerating voltage of 200 kV. For optical spectroscopy measurements, colloidal solutions of 0.02 mg/ml were prepared for each size of the QDs. Absorption spectra were obtained in standard quartz cuvettes using an HR2000 CG UV–vis–NIR spectrometer equipped with an Ocean Optics DH2000 deuterium/halogen light source.

Photocatalytic Hydrogen Evolution. Irradiation tests were performed by dispersing 50 mg of the QDs in 100 ml of aqueous 0.10 M Na2SO3 solution (as the sacrificial electron donor) in a borosilicate glass flask, followed by purging with argon and irradiating the solution mixture with a 300 W xenon arc lamp (400 mW/cm² at the flask surface as measured by an International Light IL1400BL photometer equipped with a GaAsP detector for 280–660 nm sensitivity range). The UV portion of the light (<330 nm) was removed by the borosilicate glass to avoid UV decomposition of the sacrificial electron donor. Evolved hydrogen was monitored as a function of time (Figure S2A) by a gas chromatograph equipped with a 60/80 Å molecular sieve column and thermal conductivity detector. The mass of the catalyst was verified gravimetrically after each irradiation using selective precipitation with hydrochloric acid. Hydrogen evolution rates were obtained from the hydrogen–time data, after correcting first for catalyst mass (Figure S2B) and then with regard to absorbed photons. The latter correction was performed by dividing the mass-normalized rates with the overlap integral of the absorbance spectrum of the respective QD and the emission spectrum of the Xe lamp (Figure S2C), as detailed elsewhere.43

Photoelectrochemical Measurements. CdSe QD-coated film (1.0 × 1.0 cm²) electrodes were prepared on F/SnO2 substrates (MTI Corporation, resistivity = 12–14 ohm/sq) by drop-coating QD solutions and drying in dark overnight at room temperature. F/SnO2 substrates were precleaned by consecutive sonications in acetone, methanol, and isopropyl alcohol. After the QD films were dried, silver wire was attached to the F/SnO2 substrate with carbon tape and sealed with polymer adhesive. Electrochemistry and photoelectrochemistry measurements were performed using a three-electrode cell with a quartz front window, a KCl-bridged saturated calomel electrode (SCE) as the reference electrode, a Pt coil as the counter electrode, and the CdSe QD film on F/SnO2 as the working electrode. A 50 ml volume of electrolyte made of 1.0 M KCl and 0.10 M Na2SO3 aqueous solution with a pH of 9.2 was added to the cell and degassed with N2 prior to the measurements. Na2SO3 was used as a sacrificial electron donor to resemble the condition in the hydrogen evolution tests. Cathodic potential scans were recorded from 0.4 to −1.7 V both in the dark and under chopped (front side) illumination with a scan rate of 10 mV/s using a Gamry Reference 600 potentiostat. Illumination was provided by a 300 W Xe arc lamp connected to the cell through a SiO2 fiber-optic cable (120 ± 20 mW/cm²) at the electrode as measured by an International Light IL1400BL photometer equipped with a GaAsP detector for 280–660 nm range). The potential measurements were calibrated using the standard potential of K[Fe(CN)6] (-0.358 V vs NHE).

Conflict of Interest: The authors declare no competing financial interest.

REFERENCES AND NOTES


